

# A simple Solution of Einstein's Field Equations

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## Abstract

In this short paper we present a very simple but non-trivial vacuum solution of Einstein's field equations of gravitation.

Consider the metric

$$ds^2 = (dt^2 - dx^2 - dy^2 - dz^2) - (w dt + w dz)^2, \quad (1)$$

where  $w = w(t, x, y, z)$  is a function of the Cartesian space-time coordinates  $x, y, z$  and  $t$ . Then a straightforward but longer computation<sup>1</sup> shows, that the components of the Riemann tensor  $R_{iklm}$  for this metric are given by

$$[R_{12lm}] = \frac{1}{2} \begin{bmatrix} 0 & X & Z & P \\ -X & 0 & 0 & -X \\ -Z & 0 & 0 & -Z \\ -P & X & Z & 0 \end{bmatrix} = [-R_{24lm}], \quad (2a)$$

$$[R_{13lm}] = \frac{1}{2} \begin{bmatrix} 0 & Z & Y & Q \\ -Z & 0 & 0 & -Z \\ -Y & 0 & 0 & -Y \\ -Q & Z & Y & 0 \end{bmatrix} = [-R_{34lm}], \quad (2b)$$

$$[R_{14lm}] = \frac{1}{2} \begin{bmatrix} 0 & P & Q & U \\ -P & 0 & 0 & -P \\ -Q & 0 & 0 & -Q \\ -U & P & Q & 0 \end{bmatrix}, \quad [R_{23lm}] = 0, \quad (2c)$$

where  $i, k, l, m$  range from 1 to 4 and

$$X = \frac{\partial^2 w^2}{\partial x^2}, \quad Y = \frac{\partial^2 w^2}{\partial y^2}, \quad Z = \frac{\partial^2 w^2}{\partial x \partial y}, \quad (3a)$$

$$U = \frac{\partial^2 w^2}{\partial t^2} + \frac{\partial^2 w^2}{\partial z^2} - 2 \frac{\partial^2 w^2}{\partial t \partial z}, \quad (3b)$$

$$P = \frac{\partial^2 w^2}{\partial x \partial z} - \frac{\partial^2 w^2}{\partial x \partial t}, \quad (3c)$$

$$Q = \frac{\partial^2 w^2}{\partial y \partial z} - \frac{\partial^2 w^2}{\partial y \partial t}. \quad (3d)$$

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<sup>1</sup>In fact I have computed explicit formulas for the Riemann and Ricci tensors of a particular class of metrics containing the example presented here, but these results will be described in a further paper.

All remaining components of the Riemann tensor can be obtained from its well-known symmetries [3]. Consequently, the Ricci tensor  $R_{ik}$  for the metric is

$$[R_{ik}] = \frac{1}{2} \begin{bmatrix} (w^2 - 1)U - V & P & Q & w^2 U - V \\ P & 0 & 0 & P \\ Q & 0 & 0 & Q \\ w^2 U - V & P & Q & (w^2 + 1)U - V \end{bmatrix}, \quad (4)$$

where the abbreviation

$$V = X + Y = \frac{\partial^2 w^2}{\partial x^2} + \frac{\partial^2 w^2}{\partial y^2} \quad (5)$$

has been used. The metric therefore solves the vacuum field equations [1]

$$R_{ik} = 0, \quad (6)$$

if

$$U = V = P = Q = 0. \quad (7)$$

It is quite easy to find functions which satisfy (7). For example

$$\begin{aligned} w(t) &= \sqrt{a t + b}, \\ w(t, z) &= a(t + z) + b, \\ w(t, z) &= \sqrt{a(t^2 - z^2) + b}, \\ w(t, z) &= a \exp(bt + bz) + c, \\ w(t, x, y, z) &= \sqrt{a t + b x + c y + d z + e} \end{aligned}$$

do the job. However, in all these cases we also have

$$X = Y = Z = 0, \quad (8)$$

so that the Riemann tensors vanish and the metrics are merely parametrizations of the flat Minkowski space-time. But if we choose

$$w(x, y) = \sqrt{a \ln(b x^2 + b y^2) + c}, \quad (9)$$

which is a solution of the Laplace equation  $V = 0$ , then (7) is true, but (8) does not hold. We thus have found a simple non-trivial stationary and axially symmetric solution of the field equations.

Since the metric (1), (9) does probably not have any physical relevance, we do not go into a deeper discussion of the real parameters  $a, b, c$  here. It should merely be noticed, that the metric looks like an vacuum analogue of the Gödel solution [2], but this must be discussed in detail at another place.

## References

- [1] Albert Einstein, *Die Grundlage der allgemeinen Relativitätstheorie*, Annalen der Physik 49, 769 – 822 (1916).
- [2] Kurt Gödel, *An Example of a New Type of Cosmological Solutions of Einstein's Field Equations of Gravitation*, Reviews of Modern Physics 21, 447 – 450 (1949).
- [3] Eberhard Klingbeil, *Tensorrechnung für Ingenieure*, Hochschultaschenbücher, Bibliographisches Institut, Mannheim (1966).